

# SDC for Rayleigh-Benard convection with spectral methods

December 17, 2024 | Thomas Baumann | Jülich Supercomputing Centre

# Spectral methods

## General idea

- Approximate solution via finite series expansion in some basis:  $u(x) \approx \sum_{n=0}^{N-1} u_n f_n(x)$
- Known derivative relationships of the basis allows to compute derivatives of the solution
- Converges quickly with select bases and smooth solutions

## Polynomial spectral method

- Basis functions:  $p_n(x) = x^n$
- Write vectorially:  
 $u(x) \approx \sum_{n=0}^{N-1} u_n x^n \rightsquigarrow \vec{u} = (u_0, u_1, \dots, u_{N-1})^T$
- Derivative:  $\partial_x p_n(x) = n p_{n-1}(x)$
- Pretty useless in practice!

Derivative matrix

$$\partial_x \vec{u} = \begin{pmatrix} 0 & 1 & 0 & \dots & 0 \\ 0 & 0 & 2 & \dots & 0 \\ \vdots & \ddots & \ddots & \ddots & \vdots \\ 0 & 0 & 0 & \dots & N-1 \\ 0 & 0 & 0 & \dots & 0 \end{pmatrix} \begin{pmatrix} u_0 \\ u_1 \\ \vdots \\ u_{N-2} \\ u_{N-1} \end{pmatrix}$$

# Fourier spectral method

- Basis functions:  $W_n(x) = \exp(-2\pi i x n/L)$
- Derivative:  $\partial_x W_n(x) = \frac{-2\pi i n}{L} W_n(x) \rightarrow$   
diagonal
- Compute expansion via FFT
- Works only for periodic boundary conditions (BCs)

Derivative matrix

$$\frac{-2\pi i}{L} \begin{pmatrix} 0 & & & \\ & 1 & & \\ & & \ddots & \\ & & & N \end{pmatrix}$$

# Chebyshev spectral method

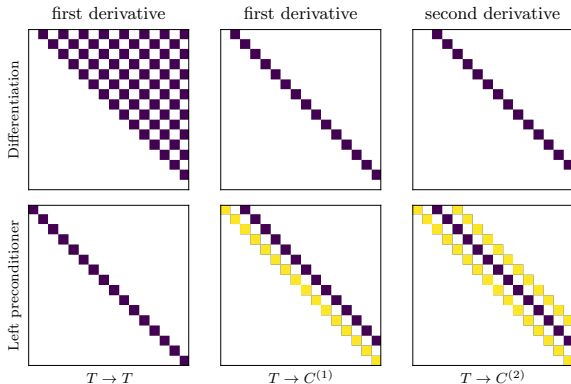
- Basis functions  $T_n(x) = \cos(n \cos^{-1}(x))$
- Bounded  $\|T_n(x)\|_\infty = 1$
- With change of variable  $x = \cos(\theta)$ :  $T_n(x) = \cos(n\theta)$ 
  - Use DCT to compute series expansion
  - Change of variables defines grid on  $x \in (-1,1)$ , points clustered at the boundary
- Use non-periodic BCs with  $\tau$ -method (see later)
- Downside: Derivative matrix is dense

$$\partial_x T_n(x) = \sum_{i=0}^{n-1} \frac{2n((n-i)\%2)}{1 + \delta_{i0}} T_i(x)$$

→ Use  $T_n$  in combination with other bases

# Ultraspherical spectral method

## Sparse preconditioned Chebychev method



- Compute expansion in Chebychev  $T$  polynomials
- Change basis while computing derivative, e.q.  $\partial_x T_n(x) = nC_{n-1}^{(1)}(x)$
- New basis: Normalized Gegenbauer polynomials  $C^{(\lambda)}$
- Render Chebychev matrices sparse with left preconditioner

# $\tau$ -method for boundary conditions

## Problem

Consider linear PDE  $L\vec{u} = \vec{a}$  with BC  $\vec{b}\vec{u} = c$  and  $N$  DoF.

$\tau$ -method: perturb PDE with  $\tau$  term in the highest mode and add BC to the system

$$\left( \begin{array}{ccc|c} L_{0,0} & \dots & L_{0,N-1} & 0 \\ \vdots & \ddots & \vdots & \vdots \\ L_{N-2,0} & \dots & L_{N-2,N-1} & 0 \\ L_{N-1,0} & \dots & L_{N-1,N-1} & 1 \\ \hline b_0 & \dots & b_{N-1} & 0 \end{array} \right) \begin{pmatrix} u_0 \\ \vdots \\ u_{N-2} \\ u_{N-1} \\ \tau \end{pmatrix} = \begin{pmatrix} a_0 \\ \vdots \\ a_{N-2} \\ a_{N-1} \\ c \end{pmatrix}$$

Get Dirichlet BCs by evaluating the elements at the boundary. E.g.  $u(x) = c$ :

$$u(x) = \sum_{n=0}^{N-1} u_n f_n(x) \rightsquigarrow b_n = f_n(x).$$

# $\tau$ -method: Example in polynomial base

## Problem

PDE:  $u_x = -1$ , BC:  $u(-1) = 1$ ,  $N = 3$ ,  $u^*(x) = -x$

## Resulting discretization

$$\left( \begin{array}{ccc|c} 0 & 1 & 0 & 0 \\ 0 & 0 & 2 & 0 \\ 0 & 0 & 0 & 1 \\ \hline 1 & -1 & 1 & 0 \end{array} \right) \begin{pmatrix} u_0 \\ u_1 \\ u_2 \\ \tau \end{pmatrix} = \begin{pmatrix} -1 \\ 0 \\ 0 \\ 1 \end{pmatrix} \rightarrow \vec{u} = \begin{pmatrix} 0 \\ -1 \\ 0 \end{pmatrix} \quad \tau = 0.$$

In practice, we don't need to compute the  $\tau$  terms and can simplify to

$$\left( \begin{array}{ccc|c} 0 & 1 & 0 & \\ 0 & 0 & 2 & \\ \hline 1 & -1 & 1 & \end{array} \right) \vec{u} = \begin{pmatrix} -1 \\ 0 \\ 1 \end{pmatrix}.$$

# $\tau$ -method: Example in polynomial base

## Problem

PDE:  $u_x = -1$ , BC:  $u(-1) = 1$ ,  $N = 3$ ,  $u^*(x) = -x$

## Resulting discretization

$$\left( \begin{array}{ccc|c} 0 & 1 & 0 & 0 \\ 0 & 0 & 2 & 0 \\ 0 & 0 & 0 & 1 \\ \hline 1 & -1 & 1 & 0 \end{array} \right) \begin{pmatrix} u_0 \\ u_1 \\ u_2 \\ \tau \end{pmatrix} = \begin{pmatrix} -1 \\ 0 \\ 0 \\ 1 \end{pmatrix} \rightarrow \vec{u} = \begin{pmatrix} 0 \\ -1 \\ 0 \end{pmatrix} \quad \tau = 0.$$

In practice, we don't need to compute the  $\tau$  terms and can simplify to

$$\left( \begin{array}{ccc|c} 0 & 1 & 0 & \\ 0 & 0 & 2 & \\ \hline 1 & -1 & 1 & \end{array} \right) \vec{u} = \begin{pmatrix} -1 \\ 0 \\ 1 \end{pmatrix}.$$



# $\tau$ -perturbations may show up in (SDC) residual

## Poisson problem in regular polynomials

PDE:  $\Delta u = 12x^2$ , BCs:  $u(-1) = -1$  and  $u(1) = 1$ ,  $u^*(x) = x^4 + x - 1$

### Properly resolved $N = 5$

$$\begin{pmatrix} 0 & 0 & 2 & 0 & 0 \\ 0 & 0 & 0 & 6 & 0 \\ 0 & 0 & 0 & 0 & 12 \\ \hline 1 & 1 & 1 & 1 & 1 \\ 1 & -1 & 1 & -1 & 1 \end{pmatrix} \vec{u}_5 = \begin{pmatrix} 0 \\ 0 \\ 12 \\ \hline 1 \\ -1 \end{pmatrix}$$

$$\vec{u}_5 = (-1, 1, 0, 0, 1)^T$$

$$r = \|\Delta u_5 - 12x^2\| = \|\Delta(x^4 + x - 1) - 12x^2\| = 0$$

### Under-resolved $N = 4$

$$\begin{pmatrix} 0 & 0 & 2 & 0 \\ 0 & 0 & 0 & 6 \\ \hline 1 & 1 & 1 & 1 \\ 1 & -1 & 1 & -1 \end{pmatrix} \vec{u}_4 = \begin{pmatrix} 0 \\ 0 \\ \hline 1 \\ -1 \end{pmatrix}$$

$$\vec{u}_4 = (0, 1, 0, 0)^T$$

$$r = \|\Delta u_4 - 12x^2\| = \|\Delta x - 12x^2\| = \|12x^2\|$$

# Rayleigh Benard convection (RBC)

## Equations

$$\begin{aligned}u_t - \nu(u_{xx} + u_{zz}) + p_x &= -uu_x - vu_z, \\v_t - \nu(v_{xx} + v_{zz}) + p_z - T &= -uv_x - vv_z, \\T_t - \kappa(T_{xx} + T_{zz}) &= -uT_x - vT_z, \\u_x + v_z &= 0.\end{aligned}$$

$$\begin{aligned}u(z = -1) &= u(z = 1) = 0 \\v(z = -1) &= v(z = 1) = 0 \\T(z = 1) &= 0, \quad T(z = -1) = 2 \\ \int_{\Omega} p &= 0\end{aligned}$$

## Discretization

- Fourier horizontally
- ultraspherical vertically

$$\Omega = [0,8) \times (-1,1)$$

Use Kronecker product to get 2D discretization from 2 1D discretizations

# IMEX Euler solver

$$M\vec{u}_t + L\vec{u} = f_{\text{nonlin}}(\vec{u})$$

$$\vec{u} = (u, v, T, p)^T$$

$$M = \text{diag}(1, 1, 1, 0)$$

$$L = \begin{pmatrix} -\partial_x^2 - \partial_z^2 & 0 & 0 & \partial_x \\ 0 & -\partial_x^2 - \partial_z^2 & -1 & \partial_z \\ 0 & 0 & -\partial_x^2 - \partial_z^2 & 0 \\ \partial_x & \partial_z & 0 & 0 \end{pmatrix}$$

$$f_{\text{nonlin}}(\vec{u}) = (uu_x + vu_z, uv_x + vv_z, uT_x + vT_z)^T$$

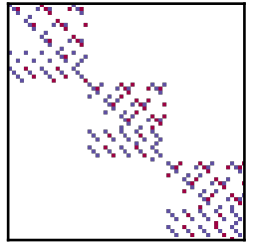
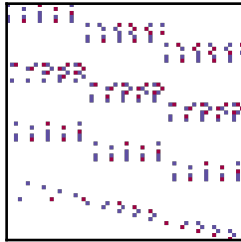
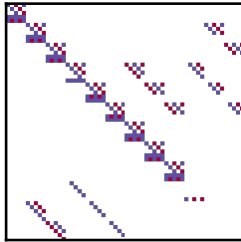
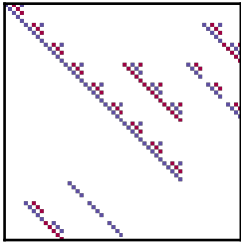
- $M$  is not invertible  $\rightarrow$  DAE
- $M + \Delta t L$  is invertible  $\rightarrow$  no worries!  
 $\rightarrow$  Need  $\tau_M = 1$ , though!
- Treat linear parts implicitly
- Treat convection explicitly

## Resulting IMEX solver

$$(M + \Delta t L) \vec{u} = M\vec{u}_0 + \Delta t f_{\text{nonlin}}(\vec{u}_0)$$

# Resulting matrix for $N = 3 \times 5$

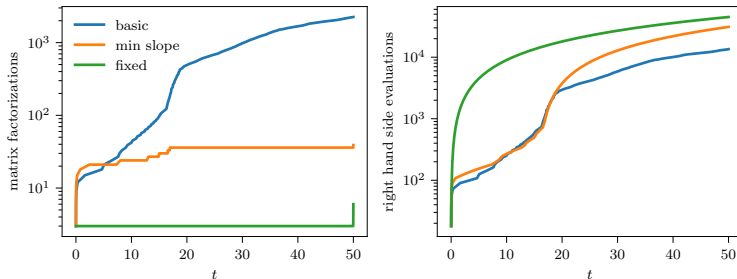
with preconditioners from Dedalus



Left to right:

- $M + L$
- + boundary conditions
- + Dirichlet recombination: Right preconditioner
- + “reverse Kronecker product”: Sort by mode, not by component → left preconditioner

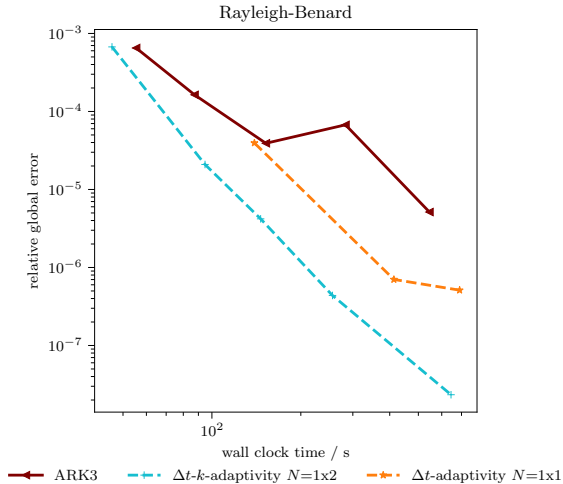
# Adaptive step size selection



## Avoiding LU factorizations via caching

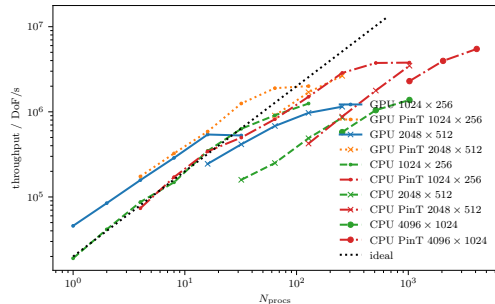
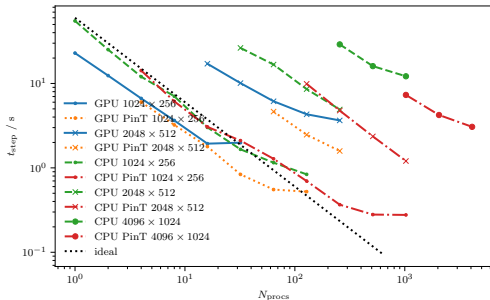
- Need to factorize when using new  $\Delta t$
- Change step size only after exceeding relative threshold
- Round step size generously to increase cache hits after restart if close to stability limit

# SDC is faster than DIRK for RBC



- Need globally stiffly accurate IMEX RKM for comparison  
→ Best I could find is order 3
- SDC at order 3 on two tasks is much faster
- Can easily increase order of SDC to obtain even greater advantage

# Parallel scaling of RBC implementation



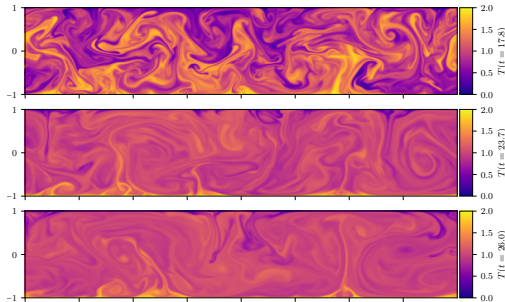
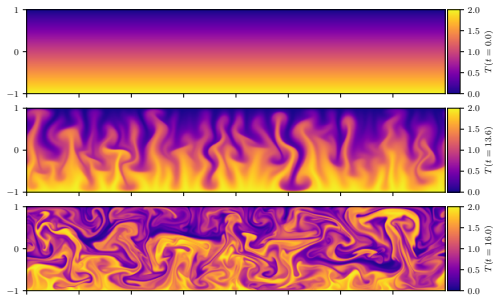
## CPUs outperform GPUs

- GPUs perform better per task
- CPUs perform better per node
- LU takes most of the time!

## Use diagonal SDC to extend scaling

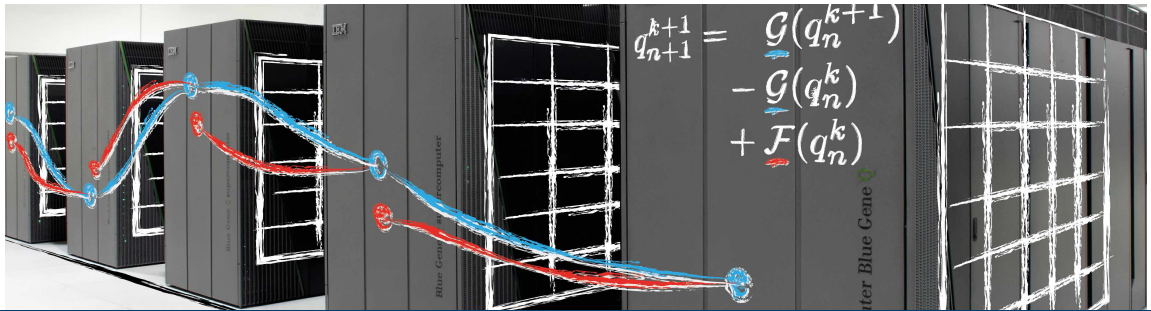
- Circumvents space-decomposition limit
- Improves strong scaling
- Enables scaling up to 4096 CPUs

# Large space-time parallel Rayleigh-Benard simulation



- Use Rayleigh number  $2 \times 10^7$  on  $N = 4096 \times 1024$  grid
- Use  $\Delta t$ - $k$ -adaptivity with four Gauß-Radau nodes for step size selection
- Use  $4 \times 1024 = 4096$  CPUs on JURECA DC (32 nodes)
- Use approx. 140k CPU hours to reach  $t = 26$  because IMEX stability limit requires  $\Delta t \approx 10^{-3}$





# SDC for Rayleigh-Benard convection with spectral methods

December 17, 2024 | Thomas Baumann | Jülich Supercomputing Centre